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2D silver-nanoplatelets metasurface for bright directional photoluminescence designed with the local Kirchhoff's law: Supplemental material

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Keywords: *metasurfaces, photoluminescence, directionality, Kirchhoff*

1 Nanoplatelets' properties

1.1 Emission and Absorption spectra of the nanoplatelets

The absorption and emission spectra of the nanoplatelets (NPLs) (in solution in hexane) are given in Figure 1. The peak emission wavelength is 605 nm. The Stokes shift, evaluated between the absorption peak (blue dotted line) and the emission peak (red dotted line) is 20 nm.

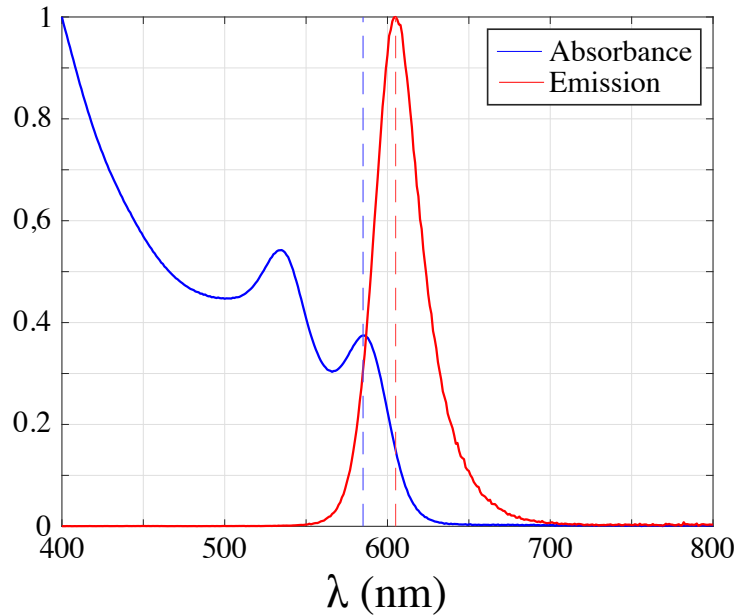


Figure 1: Normalized emission and absorption spectra of the solution of NPLs in hexane in arbitrary units.

1.2 TEM images

The TEM (transmission electron microscopy) image of the NPLs is presented in Figure 2. For TEM imaging, a drop of diluted NPLs solution in hexane is drop-casted on a copper grid covered with an amorphous carbon film. The grid is degassed overnight under secondary vacuum. A JEOL 2010F is used at 200 kV for the picture acquisition.

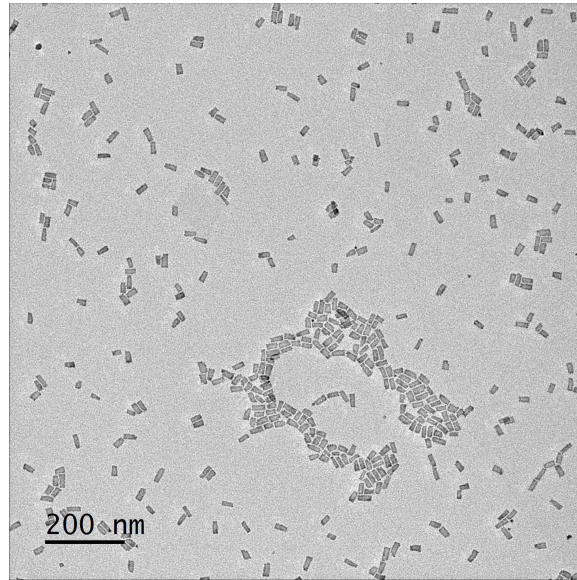


Figure 2: TEM image of the NPLs.

1.3 Refractive index of the nanoplatelets

In order to measure the refractive index of the NPLs, we fabricated a sample consisting on NPLs deposited by spin coating on top of a stack 50 nm thick layer of silver/ 1 nm thick layer of germanium/ SF10 glass substrate. Silver and germanium were deposited using electron beam evaporation. To deposit the NPLs, 200 μ L of a NPLs solution was spin coated on the metallic substrate at 500 rpm during 30 seconds with an acceleration ramp of 5 seconds. The ellipsometry measurements were performed in three steps.

1. First, the refractive index of a SF10 glass substrate was measured to serve as a reference.
2. The refractive index of silver (including 1 nm of germanium) was measured from a reference sample which was fabricated under the same conditions as the sample covered with NPLs. It thus consists in a 50 nm thick layer of silver on a 1 nm thick layer of germanium, on a SF10 glass substrate. The experimental refractive index is similar to the silver index of reference [1], as it can be seen in Figure 3, so that the germanium layer has little impact on the refractive index. The thickness of the layer was obtained by scratching it with a needle and measuring the depth of the slit by AFM. We obtained 50 ± 3 nm, in agreement with the nominal value.
3. Knowing the refractive index models of glass and silver, the refractive index of the NPLs was extracted from ellipsometry data and processed using a B-spline method (which is Kramers-Kronig consistent). The refractive index is given in Figure 4. The total thickness of the sample was obtained by scratching it with a needle and measuring the total depth of the slit by AFM. By subtracting the experimental thickness values of the silver and germanium layers obtained from the reference sample, we obtained 42 ± 5 nm thick.

In order to perform the dispersion relation computed with a complex frequency presented in Figure 5 in the main article [2], we fitted the index of the NPLs as well as the index of silver by a polynomial of degree 2: $p(\lambda) = p_1\lambda^n + p_2\lambda^{n-1} + \dots + p_n\lambda + p_{n+1}$, with the Matlab[®] function "polyfit". The fitting coefficients are given in Table 1. The comparisons between the ellipsometry measurements and the polynomial fits are presented in Figure 3 for the silver and in Figure 4 for the NPLs. A higher degree of the polynomial fits more accurately the experimental data, but the dispersion relation remains the same. Nevertheless, the absorptivity computations presented in the main article [2] are done with an interpolation of experimental index of the NPLs obtained by ellipsometry and the refractive index of silver of reference [1].

	NPLs	Ag
p_1	$0.5186 + 0.3912i$	$1.0159 - 3.7984i$
p_2	$-0.9402 - 0.7008i$	$-1.2376 + 12.7632i$
p_3	$2.1514 + 0.3220i$	$0.5036 - 2.5394i$

Table 1: Coefficients for the polynomial fit $p(\lambda) = p_1\lambda^n + p_2\lambda^{n-1} + \dots + p_n\lambda + p_{n+1}$, of degree $n = 2$ of the refractive index of the nanoplatelets (NPLs) of Figure 4 and the refractive index of silver (Ag) of Figure 3.

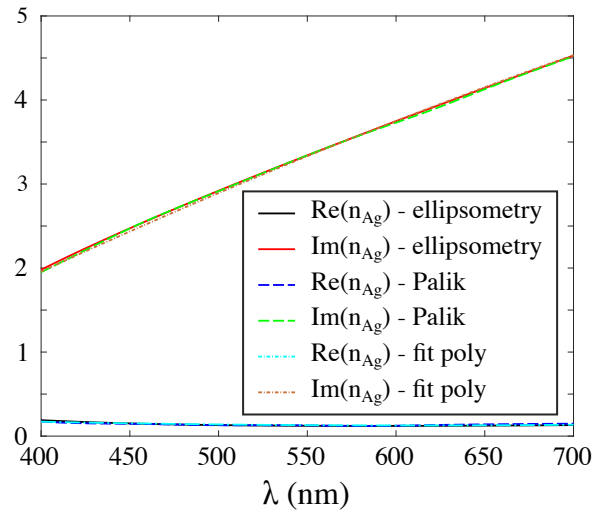


Figure 3: Refractive index of silver, measured by ellipsometry from a sample composed of 50 nm of silver on top of a 1 nm of Germanium on a SF10 glass substrate. The experimental values are compared with the refractive index of silver of reference [1], named "Palik" and the polynomial fit whose coefficients are given in Table 1.

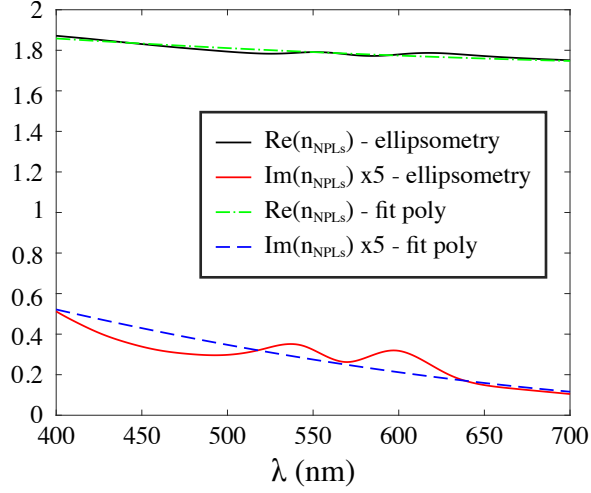


Figure 4: Refractive index of the NPLs measured by ellipsometry and comparison with the polynomial fit whose coefficients are given in Table 1, from a sample composed of 42 ± 5 nm thick layer of NPLs deposited by spin coating on top of a 50 nm of silver on a 1 nm of Germanium, on a SF10 glass substrate. For the sake of clarity, the imaginary part of the refractive index is multiplied by 5.

2 Spatial structure of the surface plasmon

This section shows the spatial structure of the mode which exists at the interface between a silver substrate and a thin layer (2 nm) of NPLs at 607.4 nm, computed with the refractive index of the NPLs measured by ellipsometry and the refractive index of silver of reference [1]. Figure 5 shows that the mode is evanescent both in metal and in air. It corresponds to a surface plasmon polariton at the interface metal/NPLs/air.

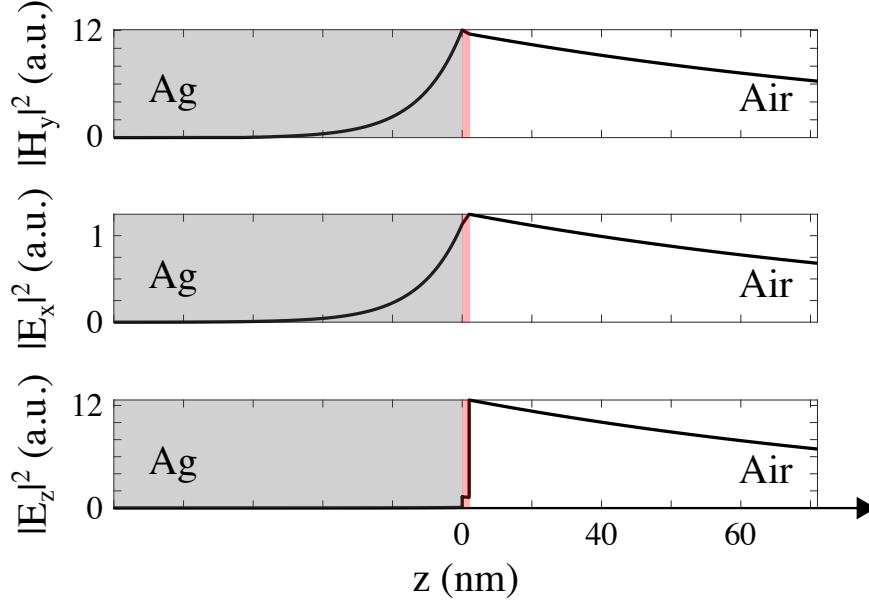


Figure 5: Spatial structure of the mode as a function of z at 607.4 nm (in arbitrary units), for a sample composed of 2 nm thick layer of NPLs (in pink hue area) on silver. At this wavelength the refractive indexes are $n_{\text{NPLs}} = 1.7844 + 0.0598i$, $n_{\text{Ag}} = 0.1260 + 3.7852i$ and $n_{\text{eff}} = 1.0419 + 0.0031i$.

3 Estimation of the beam efficiency

In this section, we present the computation of the beam efficiency at the experimental peak emission wavelength (at 607.4 nm), defined as $P_{\text{lobe}}/P_{\text{tot}}$, where $P_{\text{lobe}} = \int_0^{k_{\text{lobe}}} dP_e$ is the emitted power in the emission peak represented in red in Figure 6, with $k_{\text{lobe}} = 2.40 \mu\text{m}^{-1}$.

Since the signal is symmetrical in $\pm k_x$, we integrate over the positive axis only and multiply by two. The total power P_{tot} has been determined in two different ways:

- Experimentally, light is collected within a light cone limited by the numerical aperture of the objective (NA = 0.75), so that it is not possible to obtain the exact total power emitted between 0° and 90° . However, it is possible to estimate the beam efficiency with the total power collected, called $P_{\text{tot}}^{\text{min}} = \int_0^{k_{\text{min}}} dP_e$, represented with green dotted lines in Figure 6. We chose a value of $k_{\text{min}} = 7.36 \mu\text{m}^{-1}$ (corresponding to 45.3° at 607.4 nm), slightly lower than $k_{\text{NA}} = k_0 \text{NA}$, before the signal decreases (see Fig. 6). This estimation gives an overestimation of the beam efficiency value. We obtain a beam efficiency of 44.8 %.
- It is also possible to give a boundary value of the total emitted power emitted between 0 and 90° , by extrapolating the value of the emitted power at k_{min} for $k > k_{\text{min}}$. The integrated power is then $P_{\text{tot}}^{\text{max}} = \int_0^{k_{\text{max}}} dP_e$ and is represented with blue dotted lines in Figure 6. Thus, we obtain an underestimation of the beam efficiency value of 35 %.

We therefore estimate that the radiative efficiency lies in the range 35% and 44.8 %.

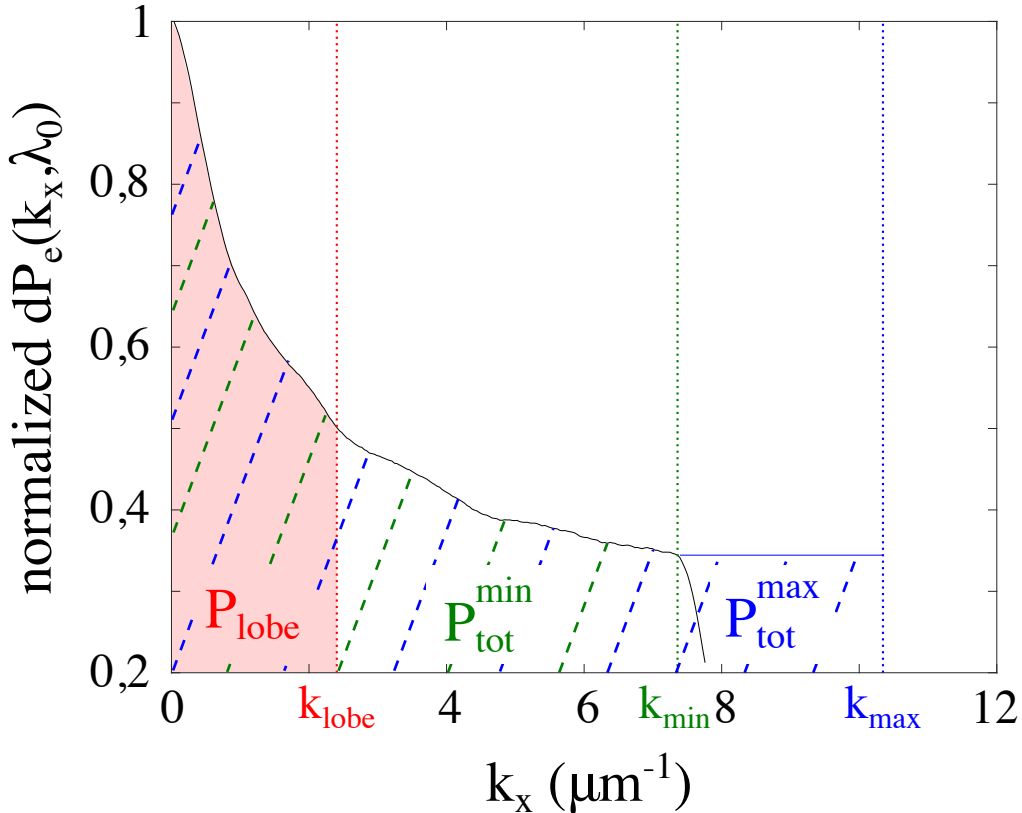


Figure 6: Normalized experimental emitted power density as a function of k_x at the peak emission wavelength, for the 2D silver metasurface covered with NPLs. Three emitted power values, each defined by different wavevector integration boundaries, can be computed: $P_{\text{lobe}} = \int_0^{k_{\text{lobe}}} dP_e$, $P_{\text{tot}}^{\text{min}} = \int_0^{k_{\text{min}}} dP_e$, $P_{\text{tot}}^{\text{max}} = \int_0^{k_{\text{max}}} dP_e$, with $k_{\text{lobe}} = 2.40 \mu\text{m}^{-1}$, $k_{\text{min}} = 7.36 \mu\text{m}^{-1}$ and $k_{\text{max}} = 10.35 \mu\text{m}^{-1}$.

4 Emission and Absorption for TE and TM polarization states

We present in Figure 7 the comparisons between normalized experimental radiation patterns and the normalized absorptivities, for Transverse Electric (TE), Transverse Magnetic (TM) polarization states, and for the total emitted power, plotted at their experimental peak emission wavelengths. Only the last case is presented in the article [2].

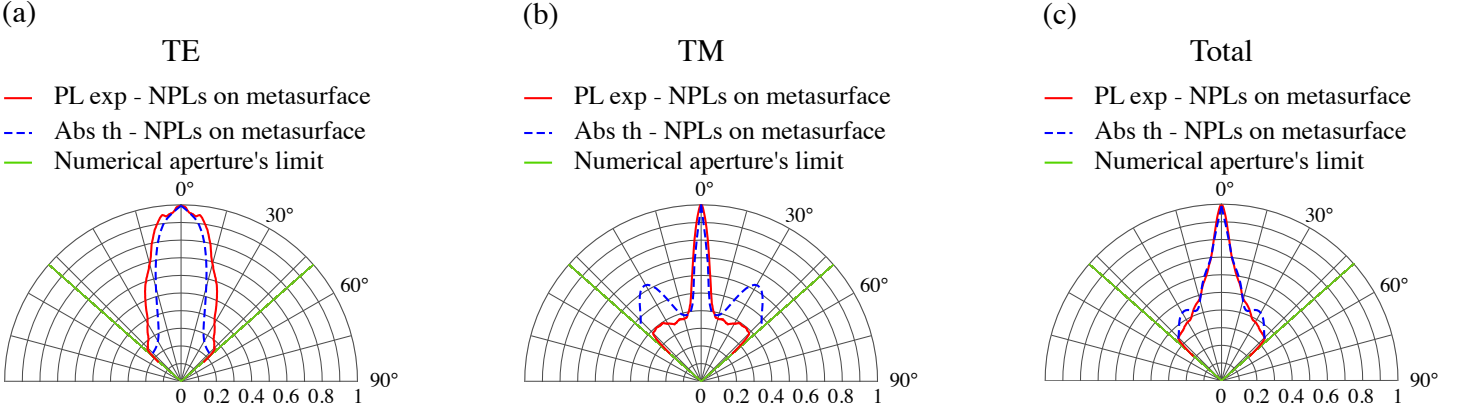


Figure 7: Normalized experimental radiation patterns at their peak emission wavelengths (red curve) and normalized theoretical absorption pattern calculated at the same peak emission wavelengths (blue dotted line), as a function of the polarization state (a): Transverse Electric (TE) at 608.4 nm, (b): Transverse Magnetic (TM) at 605.8 nm, (c): Total emission at 607.4 nm, for $h_{\text{res}} = 100$ nm, $l_{\text{res}} = 450$ nm, $p_{\text{res}} = 600$ nm, $h_{\text{top-NPLs}} = 2$ nm.

We also present in Figure 8 the comparison between the experimental radiation patterns and the normalized absorptivities for $h_{\text{top-NPLs}} = 0$ nm, that is considering there was no overfilling of the grating grooves, and showing a less good agreement.

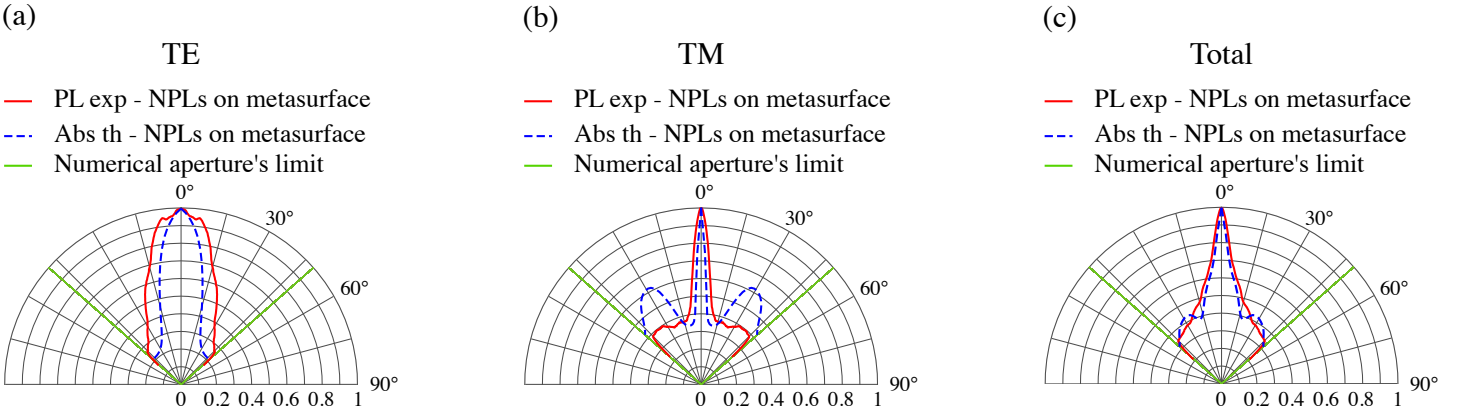


Figure 8: Normalized experimental radiation patterns at their peak emission wavelengths (red curve) and normalized theoretical absorption pattern calculated at the same peak emission wavelengths (blue dotted line), as a function of the polarization state (a): Transverse Electric (TE) at 608.4 nm, (b): Transverse Magnetic (TM) at 605.8 nm, (c): Total emission at 607.4 nm, for $h_{\text{res}} = 100$ nm, $l_{\text{res}} = 450$ nm, $p_{\text{res}} = 600$ nm, $h_{\text{top-NPLs}} = 0$ nm.

References

- [1] E. Palik, *Handbook of Optical Constants of Solids: Volume 1*. No. vol. 1, Elsevier Science, 2012.
- [2] See Article.

2D silver-nanoplatelets metasurface for bright directional photoluminescence, designed with the local Kirchhoff's law

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Keywords: *metasurfaces, photoluminescence, directionality, Kirchhoff's law*

Semiconductor colloidal nanocrystals are excellent light emitters in terms of efficiency and spectral control. Integrating them with a metasurface would pave the way to ultrathin photoluminescent devices with reduced amount of active material and performing complex functionalities such as beam shaping or polarization control. To design such a metasurface, a quantitative model of the emitted power is needed. Here, we report the design, fabrication and characterization of a ≈ 300 nm thick light-emitting device combining a plasmonic metasurface with an ensemble of nanoplatelets. The source has been designed with a new methodology based on a local form of Kirchhoff's law. The source displays record high directionality and brightness.

1 Introduction

Light emitting diodes (LEDs) have revolutionized lighting and display thanks to their efficiency and compactness. However, controlling light properties such as directionality, spectrum or polarization requires additional optical elements, such as collimators, filters or polarizers. These components reduce the optical flux and increase the size of the system. Recent technological developments in nanophotonics have demonstrated the suitability of metasurfaces as tools to control light properties using a single multifunctional flat device. Initially, metasurfaces have been designed to shape a spatially and spectrally coherent wavefront [1, 2, 3]. The building blocks of the metasurfaces behave as resonant antennas and the control on directionality is achieved via interferences of the scattered field thanks to the dephasing induced by the independent antennas.

However, this approach is unsuitable for the creation of light-emitting metasurfaces composed of many incoherent emitters integrated in the metasurface. Indeed, the emitters produce a field spatially incoherent which cannot be diffracted in a well-defined direction.

It is possible to overcome this issue by generating spatial coherence. Directional emission based on emitters deposited on a plasmonic or dielectric grating has been achieved with dye molecules [4, 5, 6, 7, 8], fluorescent glass [9], quantum wells [10, 11, 12], and more recently with semiconductor nanocrystals [13, 14, 15, 16, 17]. Restoration of a spatially coherent field was done by coupling efficiently any emitter to a delocalized mode such as a surface plasmon, a surface phonon polariton or a guided wave. A metasurface consisting of resonators can also be used if they interact so that surface lattice resonance modes exist [14]. To emit light in a controlled way, periodic perturbations are introduced so that the surface modes are converted into leaky modes.

This is the basic mechanism of directional sources of thermal radiation based on a grating ruled on a material supporting surface phonon polaritons as reported in reference [18]. A review of the state of the art of light-emitting metasurfaces can be found in the following references [19, 20].

The formalism routinely used to simulate and compute light emission by an ensemble of fluorophores embedded or in the vicinity of a metasurface consists in modelling them as point dipoles, and summing incoherently the intensities emitted by each dipole separately [19]. This technique enables to account

for the emission and polarization pattern due to the surface. However, it does not account in a self-consistent manner for the feedback of the presence of the layer of emitters on the metasurface resonances. In addition, the technique does not enable to optimize the emitted power as the dipole amplitude is unknown. In practice, the power emitted by the fluorophores depends on the pumping intensity and also on their temperature. Recently, it has been shown that the emitted power by a semiconductor can be optimized by using a reciprocity argument [21, 12]. Two issues are still pending. There is no available upper bound of the emitted power so that there is no figure of merit to characterize how far from optimum is the design in terms of brightness. The second issue is that the available models do not capture the dependence of light emission properties on temperature nor on the pumping process of the emitters.

In contrast with fluorescent sources, thermal sources are designed using Kirchhoff's law which enables to optimize simultaneously not only directionality and polarization but also emitted power by designing emissivities approaching 1 [22, 23, 24, 25, 26, 27, 28, 29, 30]. Recent works have shown that the local Kirchhoff's law can be successfully used to analyze complex photoluminescent emission situations [31, 32, 33, 34]. In this article, we show that the local form of Kirchhoff's law can be used to optimize the design of a light-emitting metasurface [35]. We extend to photoluminescence what was first achieved in the context of thermal emission, with an innovative method of designing directional photoluminescent metasurface via the local Kirchhoff's law. We then demonstrate a bright and directional photoluminescent source.

2 Experimental platform

A schematic diagram of the light-emitting metasurface is shown in Figure 1. We envision a metallic metasurface enabling the existence of surface plasmon polaritons which induce spatial coherence in the region of the interface [36]. The directionality can then be obtained by coupling the surface mode to the far field by means of a grating [18]. The whole surface is covered with emitters filling the grooves of the grating. Silver is used both as the substrate and grating material, because of its low losses in the visible spectrum. We use nanoplatelets (NPLs), because they are known to be bright emitters. NPLs are chemically synthesized core/shell semiconductor colloidal quantum wells. They consist of a 4 monolayers thick CdSe core (a monolayer is a stack of a cationic plane and an anionic plane [37]) and a 4 monolayers thick ZnS shell, with oleic acid ligands attached to the surface. This leads to a thickness of approximately 4 nm and an average lateral extension of $13 \text{ nm} \times 25 \text{ nm}$ for each NPL. Their quantum yield has been measured and is equal to 50-55 %. In the literature, a quantum yield close to unity has been obtained with core/shell CdSe/CdZnS NPLs using hot injection procedure [38, 39]. NPLs have two benefits: they display a reduced inhomogeneous broadening and they are robust to photobleaching and blinking compared to dye molecules [37]. All these properties, coupled to their low production cost, make it an interesting choice of emitter for bright light emitting devices [40]. The NPLs peak emission lies at 605 nm. Their absorption and emission spectra are given in the Supplementary Material [41].

3 Numerical design of directional photoluminescent metasurfaces with the local Kirchhoff's law

In order to optimize the design a metasurface emitting light within a narrow solid angle around the direction normal to its surface, a model of its emission is needed. Here, we use the local Kirchhoff's law [35, 42, 43]. It is based on fluctuational electrodynamics and it has been shown that it enables to model electroluminescence as well as photoluminescence [44, 31, 33] by arbitrary resonant structures containing active regions. The pumping intensity of the active regions is characterized by a photon chemical potential μ . In the specific case of electroluminescence, μ is given by the difference of the quasi-Fermi levels in the conduction and valence bands equal to eV , with e the electron charge and V the voltage bias, connected to the electrical intensity through the device by the characteristic $I(V)$ curve.

The local Kirchhoff's law [35] states that the power d^2P_e emitted by a layer of thermalized emitters at a wavelength λ in the direction \mathbf{u} and for a given polarization l can be cast in the form:

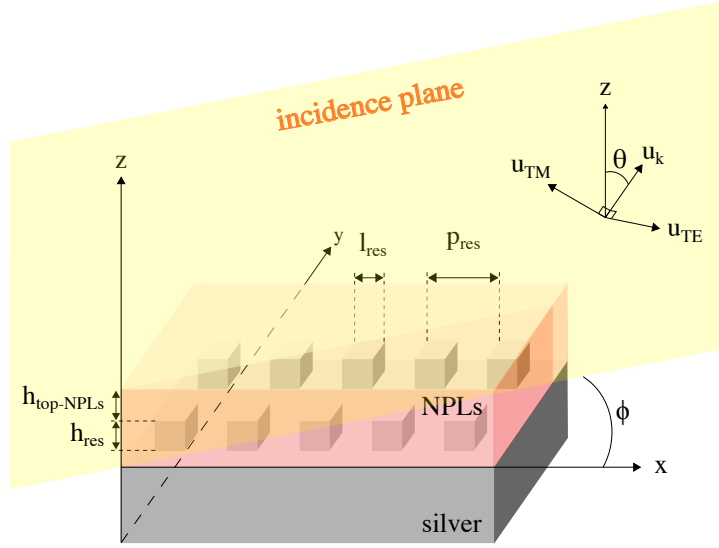


Figure 1: Schematic view of the 2D photoluminescent metasurface composed of nanoplatelets (NPLs) on top of a 2D silver grating, with $h_{\text{res}} = 100$ nm, $l_{\text{res}} = 450$ nm, $p_{\text{res}} = 600$ nm, $h_{\text{top-NPLs}} = 2$ nm. The geometric parameters are chosen in order to have a directional absorptivity (therefore reciprocally directional emissivity). The vectors \mathbf{u}_{TM} and \mathbf{u}_{TE} are unitary vectors for the Transverse Magnetic (TM) and Transverse Electric (TE) directions.

$$d^2 P_e^{(l)}(\mathbf{u}, \lambda, \mu) = d\lambda d\Omega \int_V d^3 \mathbf{r}' \alpha^{(l)}(-\mathbf{u}, \mathbf{r}', \lambda, \mu) \frac{hc^2}{\lambda^5} \frac{1}{\exp(\frac{hc}{\lambda k_B T} - \frac{\mu}{k_B T}) - 1}, \quad (1)$$

where $\alpha^{(l)}(-\mathbf{u}, \mathbf{r}', \lambda) \varphi_{\text{inc}} d^3 \mathbf{r}'$ is the absorption in a volume element $d^3 \mathbf{r}'$ of the *active material* illuminated by an incident plane wave in the $-\mathbf{u}$ direction with a given polarization l and a power flux per unit area φ_{inc} . T is the temperature and μ is the photon chemical potential. Note that Equation (1) displays explicitly a temperature dependence, as well as a pumping intensity dependence via the chemical potential μ .

For a flat source with area S , the emitted power can be expressed using the radiance $\mathcal{L}(\mathbf{u}, \lambda)$:

$$d^2 P_e^{(l)}(\mathbf{u}, \lambda) = d\lambda d\Omega S \cos \theta \mathcal{L}(\mathbf{u}, \lambda). \quad (2)$$

Assuming homogeneous temperature and chemical potential in the medium, these two equations lead to

$$\mathcal{L}^{(l)}(\mathbf{u}, \lambda, \mu, T) = \mathcal{A}^{(l)}(-\mathbf{u}, \lambda, \mu) \mathcal{L}_{\text{BB}}(\lambda, \mu, T), \quad (3)$$

where $\mathcal{A}^{(l)}(-\mathbf{u}, \lambda, \mu)$ is the absorptivity (which has an upper bound of 1), and $\mathcal{L}_{\text{BB}}(\lambda, \mu, T) = \frac{hc^2}{\lambda^5} \frac{1}{\exp(\frac{hc}{\lambda k_B T} - \frac{\mu}{k_B T}) - 1}$ is the generalized Planck's law. In an aside, we would like to point out the fact that $\mathcal{L}_{\text{BB}}(\lambda, \mu = 0, T)$ corresponds to Planck's law i.e. the blackbody radiance. Moreover, we emphasize that the absorptivity $\mathcal{A}(-\mathbf{u}, \lambda, \mu)$ must not be computed in the entire structure but only in the active region. Thus, it cannot be retrieved from the reflection and transmission factors.

Since the light directionality is given by the angular dependence of the radiance $\mathcal{L}^{(l)}(\mathbf{u}, \lambda, \mu, T)$, Equation (3) explicitly shows that it matches the angular width of the absorptivity function $\mathcal{A}^{(l)}(-\mathbf{u}, \lambda, \mu, T)$, showing that the absorptivity is the right figure of merit for light emission properties.

We now highlight the main benefits of this design methodology. The optimization of the emitted power amounts to optimize the absorptivity by the active medium. In practice, the design of a metasurface only requires a set of calculations of the absorbed power over the entire active region for a few angles and two polarization states. The technique thus requires less computation time than repeating the emission calculation by dipolar emitters averaging over all positions in a metasurface cell and over the three polarization states [19]. Furthermore, it is important to stress that Equation (1) gives an absolute expression of the emitted power. Hence, for a given chemical potential, it is seen that the emitted power has an upper bound

corresponding to an absorptivity equal to 1. Kirchhoff's law thus enables to design a metasurface which approaches a maximum absorptivity. Previous design strategies aimed at optimizing the incident field in the active region owing to reciprocity but could not quantify if the design was close from the upper bound [19, 21, 12]. In summary, Kirchhoff's law enables to design a directional source with maximum brightness.

We now turn to the implementation of this design methodology. In order to perform absorption calculations by the active region, we model the layer of NPLs by an effective homogeneous medium with a complex refractive index. This approach is known to reproduce the reflectivity of a layer of NPLs on a metal surface. It also enables to capture the strong coupling between NPLs and plasmons [45]. To do so, we do ellipsometry measurements to obtain the effective refractive index of the NPLs. The experimental data are given in the supplementary materials [41]. An accurate refractive index model for the nanoplatelets is needed in order to model the photoluminescent emitted power with Kirchhoff's law, in particular at wavelengths far from the resonance [32]. Ellipsometry data have the advantage of keeping things simple and it describes accurately the reflectivity measurements. It is thus a satisfactory model for the design of absorptivity. We also use an interpolation of the refractive index of silver of reference [46], in agreement with the ellipsometry measurements (see [41]).

To proceed, the absorptivity is computed with Rigorous Coupled-Wave Analysis (RCWA) [47, 48, 49] as a function of the angle of incidence at the peak emission wavelength of the NPLs (at 605 nm). A scan of the geometric parameters of the metasurface presented in Figure 1 is done until an absorptivity with a maximum at $\theta = 0^\circ$ is obtained, with a full-width at half-maximum (FWHM) smaller than 30° .

We now provide a guideline to optimize the absorptivity in the NPLs in order to obtain the highest possible brightness. The parameters of the metasurface have been designed so that its total absorptivity is maximal. This can be achieved if the so-called critical coupling condition is satisfied [50, 51], namely, if the radiative losses are equal to the ohmic losses in the materials. The losses of the materials are the sum of losses in the NPLs and in the metal. To enhance the absorptivity in the NPLs, we aimed at reducing the losses in the metal, and thus chose silver. As discussed in reference [52], the bare plasmonic grating must be in the overcoupled regime (radiative losses are larger than material losses) in order to ensure that material losses are equal to radiative losses when adding a layer of NPLs.

We plot the absorptivity at normal incidence in Figure 2 as a function of the wavelength (with and without the NPLs), in order to evidence the large absorptivity at critical coupling in the presence of NPLs. The solid red line corresponds to the total absorptivity (in silver and NPLs) in presence of the NPLs. At the peak emission wavelength of the NPLs, the absorptivity is close to 1. The value of absorptivity in the NPLs layer reaches 0.6 at 605 nm at normal incidence (see Figure 2, blue dotted line). These results show that the absorptivity in the NPLs, and thus the brightness, have been optimized and reaches its highest possible value.

To show that the metasurface is in the overcoupled regime, we plot the absorptivity of the bare silver grating in dark solid line with $\epsilon_{Ag} = \epsilon'_p + i\epsilon''_p$ the silver refractive index. This curve shows a broad plasmonic resonance around 825 nm whose width is given by both radiative and non-radiative losses. To check that radiative losses are dominant, we reduce the losses in the metal by imposing $\epsilon''_{Ag} = \epsilon''_p/2$. The associated absorptivity is plotted with the dark dotted line. The width of the resonance is barely modified, which proves that the resonator is over-coupled i.e. that the radiative losses dominate the losses and thus the resonance width.

Let us stress that the resonance of the metallic grating is much broader than the emitter spectral width. Hence, the same grating can be used for emitters at different wavelengths within the plasmon resonance width.

We have been able to obtain a bright and directional emission normal to the surface (see Equation (3)). With this method, we obtain a trade-off between the highest possible value of absorptivity (i.e. emissivity) at $\theta = 0^\circ$ and the directionality. For instance, increasing the thickness of the NPLs on top of the grating may lead to a higher brightness but is detrimental to the directionality, a feature that could not be identified when designing the metasurface using the average emission of dipoles. The parameters

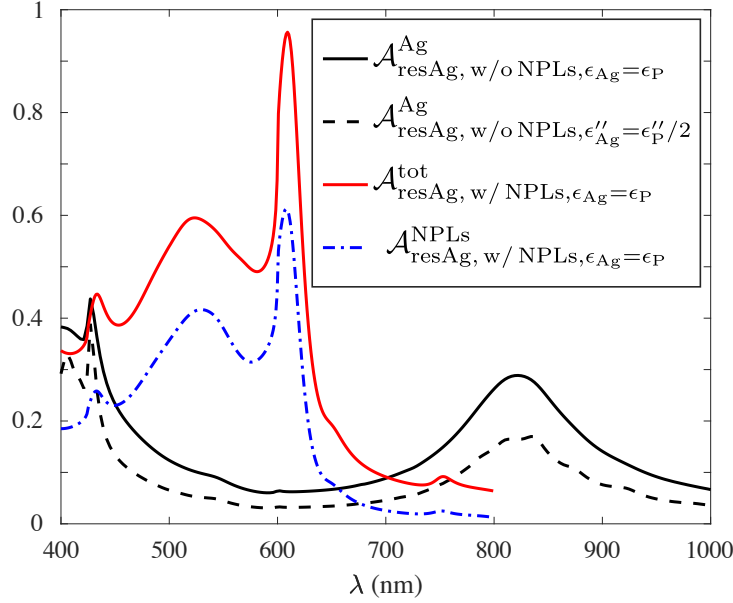


Figure 2: Absorptivity as a function of the losses in silver, for a silver grating with or without the NPLs, where the geometric parameters of the grating are given in Figure 1, at normal incidence. The dark solid line corresponds to the absorptivity in the bare silver grating, using the refractive index model of reference [46], i.e. $\epsilon_{\text{Ag}} = \epsilon'_{\text{P}} + i\epsilon''_{\text{P}}$. The dark dotted line corresponds to the absorptivity for the bare silver grating, when the losses in silver have been divided by a factor 2 i.e. $\epsilon''_{\text{Ag}} = \epsilon''_{\text{P}}/2$. The solid red line corresponds to the total absorptivity (in the silver and the NPLs) for the 2D grating with the NPLs. The blue dotted line corresponds to the absorptivity only in the NPLs for the 2D grating with the NPLs.

obtained by this procedure are equal to $h_{\text{res}} = 100$ nm, $l_{\text{res}} = 450$ nm, $p_{\text{res}} = 600$ nm and $h_{\text{top-NPLs}} = 2$ nm. Initially, we considered that the NPLs fully filled the grooves without overflowing them, so that the simulations were run with $h_{\text{top-NPLs}} = 0$ nm. However, due to the experimental spin coating deposition of the nanoplatelets (see Methods), some nanoplatelets might have been left on top of the grating. We thus adjusted the value of $h_{\text{top-NPLs}}$ and we obtained a better agreement with the experiment with $h_{\text{top-NPLs}} = 2$ nm, which corresponds to a partially filled monolayer. The numerical computations with $h_{\text{top-NPLs}} = 0$ nm are given in the supplementary.

This metasurface leads to a directional emission both in Transverse Electric (TE) and Transverse Magnetic (TM) polarizations and thus a directional total emission. Note that the design is performed using the total emitted power, adding the TE and TM polarizations. We show both polarizations in the supplementary [41].

4 Results

The sample was fabricated using electronic lithography on a silicon substrate and the layer of NPLs has been deposited by spin coating, as detailed in the Methods section. A SEM image of the grating before deposition of the nanoplatelets is presented in Figure 3. The width of the pillars has been evaluated from the SEM images, and it varies from 480 nm in the bottom to 450 nm in the top. In the numerical simulations we considered straight pillars with a uniform width of 450 nm.

The sample is pumped using the 436 nm line of a mercury vapor lamp. The light is focused on the sample by an objective with a numerical aperture of 0.75, probing an area whose diameter is around 110 μm large. The photoluminescence is then collected through the same objective and the back focal plane of the microscope objective is imaged onto the slit of a spectrometer. A Fourier plane image as well as energy-momentum figures of photoluminescence from the optically pumped metasurfaces are shown in Figure 4a,b. The emission is highly directional with a FWHM about $\pm 26.8^\circ$ at the peak wavelength (represented by the horizontal dotted line in Figure 4b).

An experimental comparison with a sample without the grating (that is with a flat, uniform and ho-

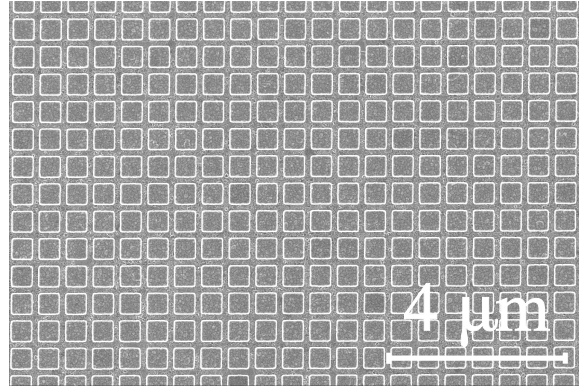


Figure 3: Scanning electron microscopy (SEM) image of a 2D grating of silver, before nanoplatelets deposition.

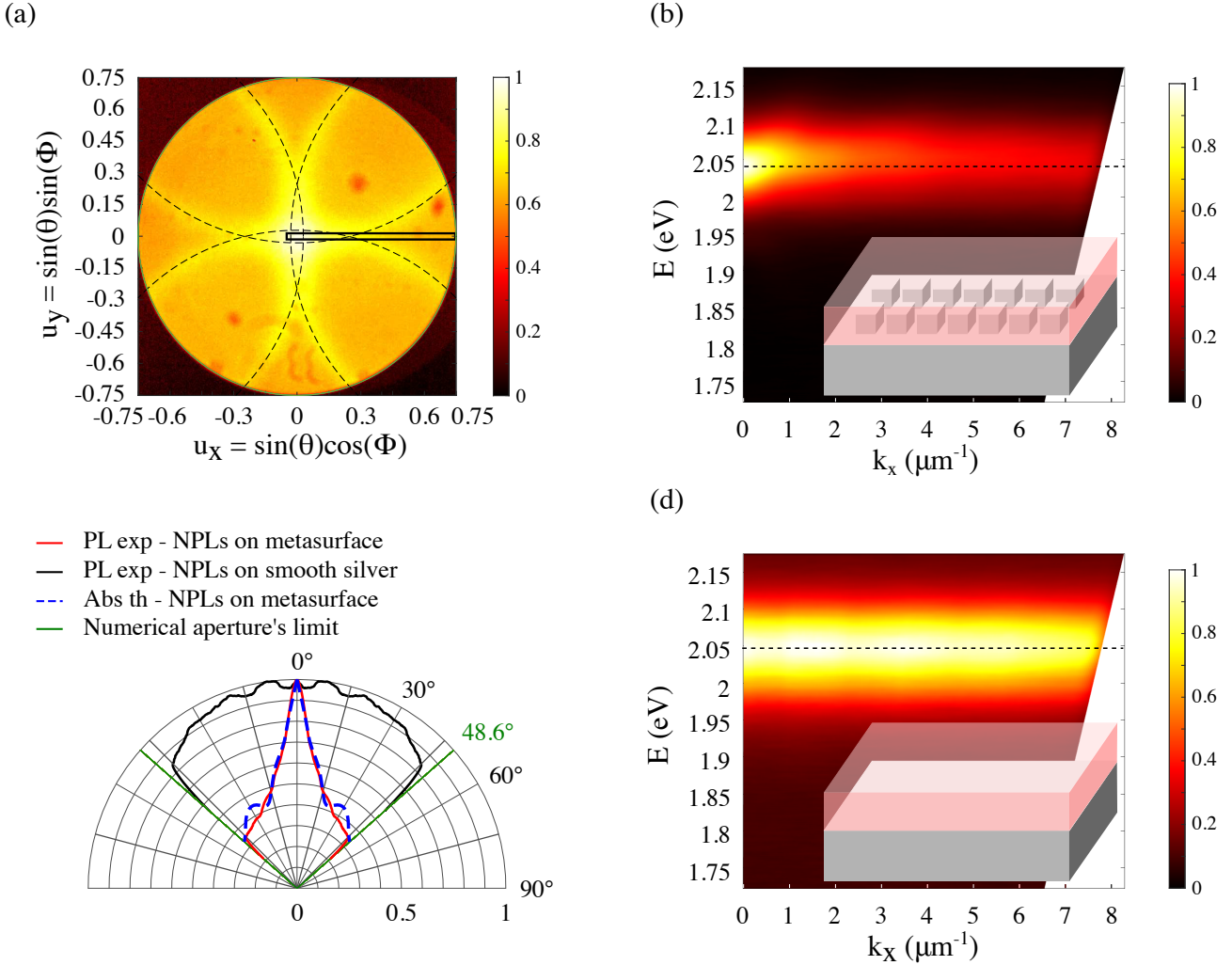


Figure 4: Experimental radiation pattern. (a) Normalized Fourier image of the sample. The dotted lines corresponds to the theoretical isofrequency projection on the (k_x, k_y) plane of the dispersion relation of a surface plasmon polariton existing in a 2 nm thick NPLs layer deposited on silver, after diffraction by a grating of $p_{\text{res}} = 600$ nm, at 607.4 nm (see Figure 5c). The black rectangle corresponds to the slit position at the entrance of the spectrometer. (b) Photoluminescence intensity of a sample consisting of NPLs on a 2D metasurface as a function of the energy and the in-plane wave vector k_x for NPLs on the 2D metasurface, (c) Normalized experimental radiation patterns for NPLs on the 2D metasurface (red curve) or on smooth silver (black line) at their experimental peak wavelength (in dotted lines in Figure b,d) and normalized theoretical absorption pattern (blue dotted line) calculated at the experimental peak emission wavelength (at 607.4 nm, black dotted line in Figure b). The green lines correspond to the numerical aperture's limit ($\text{NA} = 0.75$). (d) Photoluminescence as a function of the photon energy and the in-plane wave vector k_x for NPLs on smooth silver.

mogeneous emitter layer with no structuration) is given in Figure 4c,d. In this latter case, the emission is quasi isotropic, as expected for an ensemble of point-like emitters (black line in Figure 4c). The comparison between the experimental radiation pattern (red line in Figure 4c) and the theoretical absorptivity pattern (blue dotted line in Figure 4c) shows a good agreement between experiment and theory which validates the method of design.

To interpret the far-field emission of Figure 4a, we simply assume that the coupling is mediated by the surface plasmon propagating at the interface silver/NPLs/air. Indeed, the size of the pillars is much larger than the width of the thin grooves (Figure 3). Thus, the sample can be seen to first approximation as a translational invariant system, composed of NPLs on top of silver, which is perturbed by the presence of the grooves. Therefore, we can plot the dispersion relation of the surface plasmons propagating at the interface between a 2 nm thick layer of NPLs and a silver substrate, see Figure 5a. It has been obtained by searching numerically the poles of the reflectivity of the system, with a complex frequency (or energy E (eV)) and a real wavevector k_{\parallel} [53, 54]. To do so, we fitted the refractive indexes of the NPLs and silver by a polynome (see [41]). At the experimental peak emission wavelength (607.4 nm), the wavevector is $k_{\parallel}^{607.4} = 10.78 \mu\text{m}^{-1}$. In the plane (k_x, k_y) , the corresponding dispersion relation is a circle (Figure 5b).

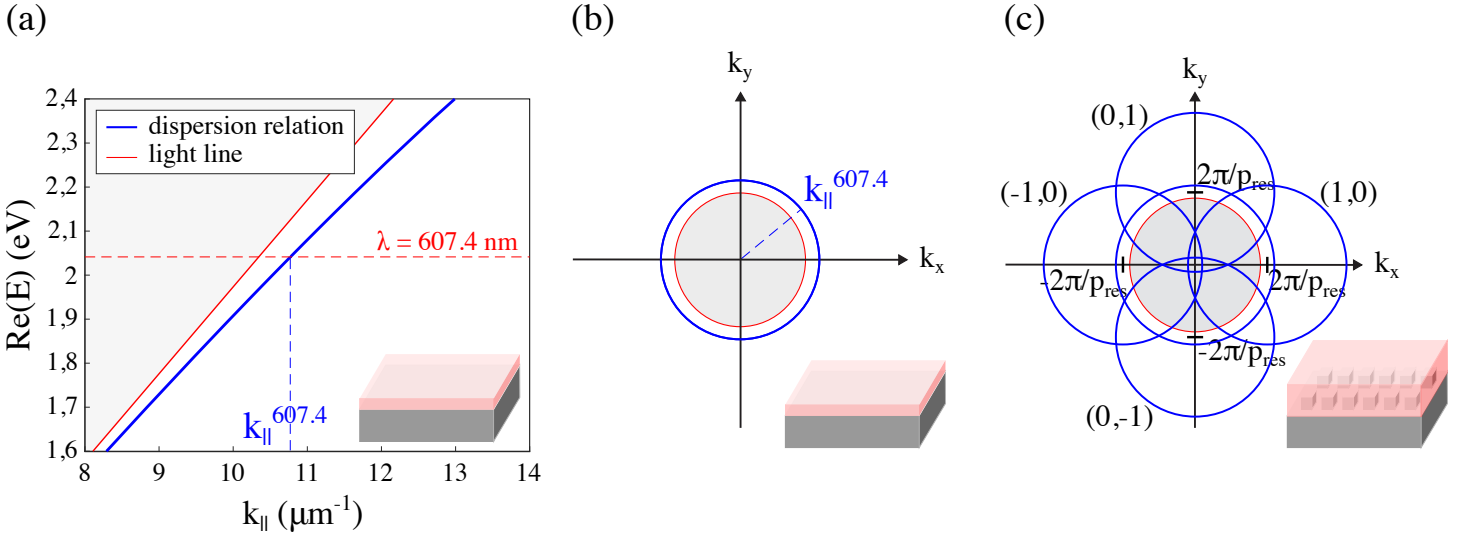


Figure 5: (a) Dispersion relation of a surface plasmon at the interface between a silver substrate and a 2 nm thick layer of NPLs. (b) Schematic diagram of its isofrequency projection on the (k_x, k_y) plane, at $\lambda = 607.4$ nm. (c) Schematic diagram of the isofrequency contour of the surface plasmon dispersion relation on the 2D metasurface with a periodicity p_{res} .

We now take into account the two-dimensional periodicity of the metasurface, by periodizing the initial dispersion relation in the momentum plane in both k_x and k_y direction (see Figure 5c). If we superimpose the circles of Figure 5c on the Fourier image of Figure 4a (normalizing the radius and the circles' center's positions by $k_0 = 2\pi/\lambda$, with $\lambda = 607.4$ nm) we show that we can reproduce accurately the observed emission pattern along the arcs of circles. In summary, the emission is enhanced by coupling to leaky surface plasmons. The choice of the periodicity is critical to ensure a peak at normal incidence.

To characterize the directionality of the source, we use a figure of merit called "beam efficiency", defined in reference [11] as the fraction of total light intensity collected within one FWHM along the k_x axis at the peak emission frequency. Since the collected light is limited by the numerical aperture, we cannot measure the total emitted intensity. We derive an upper bound and a lower bound by assuming that the emitted power into angles larger than 45.3° is either null or equal to its value at 45.3° . We obtain an efficiency in the range 35% and 44.8 %, larger than previous reports (see Supplementary [41]).

Note that the enhancement due to the metasurface enables to achieve large brightness with a small amount of active material. At the experimental peak emission wavelength, the absorptivity is 0.61. To reach the same absorptivity with a nanoplatelets layer on a glass substrate at normal incidence, the thickness layer should be 850 nm, corresponding to a volume 18.6 times larger.

5 Conclusion

In summary, we have shown that the extension of Kirchhoff’s law to photoluminescent metasurfaces provides an accurate model of the emission. To our knowledge, it is the first time that this law has been used to design a photoluminescent metasurface with specific emission properties. Here, we have shown that it is possible to design a bright metasurface emitting within a narrow spectral range and within a narrow cone with half-angle below 15° , taking advantage of surface plasmons. Such a device has been realized for the first time with nanoplatelets in the visible range. We have obtained a beam efficiency on the order of 40% as well as an absorptivity of 0.61. The same approach could be used to design photoluminescent metasurfaces controlling the emitted polarization as reported recently in the context of thermal emission [25]. The method could also be applied to optimize dielectric photoluminescent metasurfaces [21, 55, 56] which could enable sources with narrower emission cones, as well as electroluminescent metasurfaces [57, 17, 16].

6 Methods

Lift-off process for the silver metasurface:

The silver grating was fabricated by lift-off process. First, a 1-nm-thick germanium layer was deposited on a silicon substrate in order to facilitate the wetting on the substrate, followed by a 200-nm-thick layer of silver deposition, by electron beam evaporation (at 0.1 nm/s). A PMMA A4 e-beam resist was then spin coated on the sample (4000 rpm during 60 seconds with an acceleration of 2000 rpm/s) and baked at 180°C for 2 minutes. E-beam lithography was processed at 100 keV with a dose of $800\ \mu\text{C}/\text{cm}^2$ and a 5 nA current on a Raith EBPG5200 machine. The sample was then developed in MIBK:IPA (1:3) for 1 minute at 20° , rinsed with IPA (for 30 seconds) and dried with nitrogen. A 100-nm-thick layer of silver was then deposited by electron beam evaporation (at 0.1 nm/s) and the final $600 \times 600\ \mu\text{m}^2$ structure was obtained by immersing the sample the whole night in butanone. The sample was then rinsed with IPA and dried with nitrogen.

Spin coating of nanoplatelets:

The nanoplatelet solution was fabricated by Corentin Dabard and Sandrine Ithurria (Laboratoire de Physique et d’Etude des Matériaux, PSL Research University, CNRS UMR 8213, UPMC Sorbonne Université, ESPCI Paris, 10 rue Vauquelin, 75005 Paris, France). The nanoplatelets were dispersed in Toluene (with a concentration $C \approx 3.8\ \text{mg}\cdot\text{mL}^{-1}$). 200 μL of nanoplatelet solution was spincoated on the metallic grating at 500 rpm during 30 seconds with an acceleration ramp of 5 seconds. When the nanoplatelets are deposited on top of a flat silver substrate with these spin coating parameters, it leads to a layer of thickness $42 \pm 5\ \text{nm}$ thick. The thickness has been measured by doing a scratch on the sample and measuring the height of the slit by AFM [41]. When depositing the nanoplatelets on the grating sample, the nanoplatelets fill the grooves since the walls of the grating are much higher than the thickness of the nanoplatelets layer on a flat substrate. We considered that very few nanoplatelets remained on top of the grating, allowing in our simulation the parameter $h_{\text{top-NPLs}}$ to be non-zero and leading to a value of 2 nm to optimize the comparison with data.

Supporting Information

Supporting Information is available.

Acknowledgements

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References

- [1] A. Li, S. Singh, and D. Sievenpiper, “Metasurfaces and their applications,” *Nanophotonics*, vol. 7, no. 6, pp. 989–1011, 2018.
- [2] N. Yu and F. Capasso, “Flat optics with designer metasurfaces,” *Nature Materials*, vol. 13, no. 2, pp. 139–150, 2014.
- [3] N. Yu, P. Genevet, M. A. Kats, F. Aieta, J.-P. Tetienne, F. Capasso, and Z. Gaburro, “Light propagation with phase discontinuities: Generalized laws of reflection and refraction,” *Science*, vol. 334, no. 6054, pp. 333–337, 2011.
- [4] G. Vecchi, V. Giannini, and J. Gómez Rivas, “Shaping the fluorescent emission by lattice resonances in plasmonic crystals of nanoantennas,” *Phys. Rev. Lett.*, vol. 102, p. 146807, Apr 2009.
- [5] J. M. Steele, I. Gagnidze, and S. M. Wiele, “Efficient extraction of fluorescence emission utilizing multiple surface plasmon modes from gold wire gratings,” *Plasmonics*, vol. 5, no. 3, pp. 319–324, 2010.
- [6] G. Lozano, D. J. Louwers, S. Rodríguez, S. Murai, O. T. A. Jansen, M. A. Verschuuren, and J. Gómez Rivas, “Plasmonics for solid-state lighting: enhanced excitation and directional emission of highly efficient light sources,” *Light: Science & Applications*, vol. 2, p. e66, 2013.
- [7] G. Lozano, G. Grzela, M. A. Verschuuren, M. Ramezani, and J. G. Rivas, “Tailor-made directional emission in nanoimprinted plasmonic-based light-emitting devices,” *Nanoscale*, vol. 6, pp. 9223–9229, 2014.
- [8] S. Wang, Q. Le-Van, T. Peyronel, M. Ramezani, N. Van Hoof, T. G. Tiecke, and J. Gómez Rivas, “Plasmonic nanoantenna arrays as efficient etendue reducers for optical detection,” *ACS Photonics*, vol. 5, no. 6, pp. 2478–2485, 2018.
- [9] A. Vaskin, J. Bohn, K. E. Chong, T. Bucher, M. Zilk, D.-Y. Choi, D. N. Neshev, Y. S. Kivshar, T. Pertsch, and I. Staude, “Directional and spectral shaping of light emission with mie-resonant silicon nanoantenna arrays,” *ACS Photonics*, vol. 5, no. 4, pp. 1359–1364, 2018.
- [10] J. DiMaria, E. Dimakis, T. D. Moustakas, and R. Paiella, “Plasmonic off-axis unidirectional beaming of quantum-well luminescence,” *Applied Physics Letters*, vol. 103, p. 251108, 12 2013.
- [11] P. P. Iyer, R. A. DeCrescent, Y. Mohtashami, G. Lheureux, N. A. Butakov, A. Alhassan, C. Weisbuch, S. Nakamura, S. P. DenBaars, and J. A. Schuller *Nature Photonics*, vol. 14, no. 9, pp. 543–548, 2020.
- [12] L. Heki, Y. Mohtashami, R. A. DeCrescent, A. Alhassan, S. Nakamura, S. P. DenBaars, and J. A. Schuller, “Designing highly directional luminescent phased-array metasurfaces with reciprocity-based simulations,” *ACS Omega*, vol. 7, no. 26, pp. 22477–22483, 2022.
- [13] S. R. K. Rodriguez, G. Lozano, M. A. Verschuuren, R. Gómez, K. Lambert, B. De Geyter, A. Hassinen, D. Van Thourhout, Z. Hens, and J. Gómez Rivas, “Quantum rod emission coupled to plasmonic lattice resonances: A collective directional source of polarized light,” *Applied Physics Letters*, vol. 100, no. 11, p. 111103, 2012.
- [14] R. Guo, S. Derom, A. I. Väkeväinen, R. J. A. van Dijk-Moes, P. Liljeroth, D. Vanmaekelbergh, and P. Törmä, “Controlling quantum dot emission by plasmonic nanoarrays,” *Opt. Express*, vol. 23, pp. 28206–28215, Nov 2015.
- [15] X. Wang, Y. Li, R. Toufanian, L. C. Kogos, A. M. Dennis, and R. Paiella, “Geometrically tunable beamed light emission from a quantum-dot ensemble near a gradient metasurface,” *Advanced Optical Materials*, vol. 8, no. 8, p. 1901951, 2020.

- [16] C. Dabard, E. Bossavit, T. H. Dang, N. Ledos, M. Cavallo, A. Khalili, H. Zhang, R. Alchaar, G. Patriarche, A. Vasanelli, B. T. Diroll, A. Degiron, E. Lhuillier, and S. Ithurria, “Electroluminescence and plasmon-assisted directional photoluminescence from 2d hgte nanoplatelets,” *The Journal of Physical Chemistry C*, vol. 127, no. 30, pp. 14847–14855, 2023.
- [17] E. Bossavit, T. H. Dang, P. He, M. Cavallo, A. Khalili, C. Dabard, H. Zhang, D. Gacemi, M. G. Silly, C. Abadie, B. Gallas, D. Pierucci, Y. Todorov, C. Sirtori, B. T. Diroll, A. Degiron, E. Lhuillier, and A. Vasanelli, “Plasmon-assisted directional infrared photoluminescence of hgte nanocrystals,” *Advanced Optical Materials*, vol. n/a, no. n/a, p. 2300863, 2023.
- [18] J.-J. Greffet, R. Carminati, K. Joulain, J.-P. Mulet, S. Mainguy, and Y. Chen, “Coherent emission of light by thermal sources,” *Nature*, vol. 416, p. 61, 2002.
- [19] A. Vaskin, R. Kolkowski, A. F. Koenderink, and I. Staude, “Light-emitting metasurfaces,” *Nanophotonics*, vol. 8, no. 7, pp. 1151 – 1198, 2019.
- [20] G. Lozano, S. R. K. Rodriguez, M. A. Verschuuren, and J. Gómez Rivas, “Metallic nanostructures for efficient LED lighting,” *Light: Science & Applications*, vol. 5, no. 6, pp. e16080–e16080, 2016.
- [21] S. Liu, A. Vaskin, S. Addamane, B. Leung, M.-C. Tsai, Y. Yang, P. P. Vabishchevich, G. A. Keeler, G. Wang, X. He, Y. Kim, N. F. Hartmann, H. Htoon, S. K. Doorn, M. Zilk, T. Pertsch, G. Balakrishnan, M. B. Sinclair, I. Staude, and I. Brener, “Light-emitting metasurfaces: Simultaneous control of spontaneous emission and far-field radiation,” *Nano Letters*, vol. 18, no. 11, pp. 6906–6914, 2018. PMID: 30339762.
- [22] D. Costantini, A. Lefebvre, A.-L. Coutrot, I. Moldovan-Doyen, J.-P. Hugonin, S. Boutami, F. Marquier, H. Benisty, and J.-J. Greffet, “Plasmonic metasurface for directional and frequency-selective thermal emission,” *Phys. Rev. Appl.*, vol. 4, p. 014023, Jul 2015.
- [23] L. Wojszvzyk, A. Nguyen, A.-L. Coutrot, C. Zhang, B. Vest, and J.-J. Greffet, “An incandescent metasurface for quasimonochromatic polarized mid-wave infrared emission modulated beyond 10 mhz,” *Nature Communications*, vol. 12, no. 1, p. 1492, 2021.
- [24] A. Nguyen and J.-J. Greffet, “Efficiency optimization of mid-infrared incandescent sources with time-varying temperature,” *Opt. Mater. Express*, vol. 12, pp. 225–239, Jan 2022.
- [25] A. Nguyen, J.-P. Hugonin, A.-L. Coutrot, E. Garcia-Caurel, B. Vest, and J.-J. Greffet, “Large circular dichroism in the emission from an incandescent metasurface,” *Optica*, vol. 10, pp. 232–238, Feb 2023.
- [26] C. M. Watts, X. Liu, and W. J. Padilla, “Metamaterial electromagnetic wave absorbers,” *Advanced Materials*, vol. 24, no. 23, pp. OP98–OP120, 2012.
- [27] D. G. Baranov, Y. Xiao, I. A. Nechepurenko, A. Krasnok, A. Alù, and M. A. Kats, “Nanophotonic engineering of far-field thermal emitters,” *Nature Materials*, vol. 18, no. 9, pp. 920–930, 2019.
- [28] A. C. Overvig, S. A. Mann, and A. Alù, “Thermal metasurfaces: Complete emission control by combining local and nonlocal light-matter interactions,” *Phys. Rev. X*, vol. 11, p. 021050, Jun 2021.
- [29] C. Arnold, F. Marquier, M. Garin, F. Pardo, S. Collin, N. Bardou, J.-L. Pelouard, and J.-J. Greffet, “Coherent thermal infrared emission by two-dimensional silicon carbide gratings,” *Phys. Rev. B*, vol. 86, p. 035316, Jul 2012.
- [30] X. Wang, T. Sentz, S. Bharadwaj, S. K. Ray, Y. Wang, D. Jiao, L. Qi, and Z. Jacob, “Observation of nonvanishing optical helicity in thermal radiation from symmetry-broken metasurfaces,” *Science Advances*, vol. 9, no. 4, p. eade4203, 2023.
- [31] H. Monin, A. Loirette-Pelous, E. De Leo, A. A. Rossinelli, F. Prins, D. J. Norris, E. Bailly, J.-P. Hugonin, B. Vest, and J.-J. Greffet, “Controlling light emission by a thermalized ensemble of colloidal quantum dots with a metasurface,” *Opt. Express*, vol. 31, pp. 4851–4861, Jan 2023.

- [32] E. Bailly, K. Chevrier, C. Perez de la Vega, J.-P. Hugonin, Y. De Wilde, V. Krachmalnicoff, B. Vest, and J.-J. Greffet, “Method to measure the refractive index for photoluminescence modelling,” *Opt. Mater. Express*, vol. 12, pp. 2772–2781, Jul 2022.
- [33] C. R. Pérez de la Vega, E. Bailly, K. Chevrier, B. Vest, J.-P. Hugonin, A. Bard, A. Gassenq, C. Symonds, J.-M. Benoit, J. Bellessa, J.-J. Greffet, Y. De Wilde, and V. Krachmalnicoff, “Plasmon-mediated energy transfer between two systems out of equilibrium,” *ACS Photonics*, vol. 10, no. 4, pp. 1169–1176, 2023.
- [34] A. Caillas, S. Suffit, P. Filloux, E. Lhuillier, and A. Degiron, “Identification of two regimes of carrier thermalization in pbs nanocrystal assemblies,” *The Journal of Physical Chemistry Letters*, vol. 12, no. 21, pp. 5123–5131, 2021. PMID: 34029086.
- [35] J.-J. Greffet, P. Bouchon, G. Brucoli, and F. Marquier, “Light emission by nonequilibrium bodies: local kirchhoff law,” *Physical Review X*, vol. 8, no. 2, p. 021008, 2018.
- [36] E. Bailly, J.-P. Hugonin, B. Vest, and J.-J. Greffet, “Spatial coherence of light emitted by thermalized ensembles of emitters coupled to surface waves,” *Phys. Rev. Res.*, vol. 3, p. L032040, Aug 2021.
- [37] B. T. Diroll, B. Guzelturk, H. Po, C. Dabard, N. Fu, L. Makke, E. Lhuillier, and S. Ithurria, “2d ii–vi semiconductor nanoplatelets: From material synthesis to optoelectronic integration,” *Chemical Reviews*, vol. 123, no. 7, pp. 3543–3624, 2023. PMID: 36724544.
- [38] Y. Altintas, U. Quliyeva, K. Gungor, O. Erdem, Y. Kelestemur, E. Mutlugun, M. V. Kovalenko, and H. V. Demir, “Highly stable, near-unity efficiency atomically flat semiconductor nanocrystals of cdse/zns hetero-nanoplatelets enabled by zns-shell hot-injection growth,” *Small*, vol. 15, no. 8, p. 1804854, 2019.
- [39] A. A. Rossinelli, A. Riedinger, P. Marqués-Gallego, P. N. Knüsel, F. V. Antolinez, and D. J. Norris, “High-temperature growth of thick-shell cdse/cds core/shell nanoplatelets,” *Chem. Commun.*, vol. 53, pp. 9938–9941, 2017.
- [40] H. D. Baruj, I. Bozkaya, B. Canimkurbey, A. T. Isik, F. Shabani, S. Delikanli, S. Shendre, O. Erdem, F. Isik, and H. V. Demir, “Highly-directional, highly-efficient solution-processed light-emitting diodes of all-face-down oriented colloidal quantum well self-assembly,” *Small*, vol. 19, no. 29, p. 2206582, 2023.
- [41] See Supplemental Material.
- [42] S. M. Rytov, Y. A. Kravtsov, and V. I. Tatarskii, *Principles of statistical radiophysics. 3. Elements of random fields*, vol. 3. Berlin: Springer, 1989.
- [43] P. Wurfel, “The chemical potential of radiation,” *Journal of Physics C: Solid State Physics*, vol. 15, no. 18, p. 3967, 1982.
- [44] H. Wang, A. Aassime, X. Le Roux, N. J. Schilder, J.-J. Greffet, and A. Degiron, “Revisiting the role of metallic antennas to control light emission by lead salt nanocrystal assemblies,” *Physical Review Applied*, vol. 10, no. 3, p. 034042, 2018.
- [45] I. Shlesinger, H. Monin, J. Moreau, J.-P. Hugonin, M. Dufour, S. Ithurria, B. Vest, and J.-J. Greffet, “Strong Coupling of Nanoplatelets and Surface Plasmons on a Gold Surface,” *ACS Photonics*, vol. 6, no. 11, pp. 2643–2648, 2019.
- [46] E. Palik, *Handbook of Optical Constants of Solids: Volume 1*. No. vol. 1, Elsevier Science, 2012.
- [47] M. G. Moharam and T. K. Gaylord, “Rigorous coupled-wave analysis of planar-grating diffraction,” *J. Opt. Soc. Am.*, vol. 71, pp. 811–818, Jul 1981.

- [48] L. Li, “New formulation of the fourier modal method for crossed surface-relief gratings,” *J. Opt. Soc. Am. A*, vol. 14, pp. 2758–2767, Oct 1997.
- [49] J.-P. Hugonin and P. Lalanne, “Reticolo software for grating analysis,” 2021.
- [50] H. Haus, *Waves and Fields in Optoelectronics*. Prentice-Hall series in solid state physical electronics, Prentice-Hall, 1984.
- [51] H. Benisty, J. Greffet, and P. Lalanne, *Introduction to Nanophotonics*. Oxford Graduate Texts, Oxford University Press, 2022.
- [52] L. Paggi, A. Fabas, H. El Ouazzani, J.-P. Hugonin, N. Fayard, N. Bardou, C. Dupuis, J.-J. Greffet, and P. Bouchon, “Over-coupled resonator for broadband surface enhanced infrared absorption (seira),” *Nature Communications*, vol. 14, no. 1, p. 4814, 2023.
- [53] A. Archambault, T. V. Teperik, F. Marquier, and J.-J. Greffet, “Surface plasmon fourier optics,” *Phys. Rev. B*, vol. 79, p. 195414, May 2009.
- [54] H. Benisty, J.-J. Greffet, and P. Lalanne, *Introduction to Nanophotonics*. Oxford Graduate Texts, Oxford University Press, 2022.
- [55] M. Decker and I. Staude, “Resonant dielectric nanostructures: a low-loss platform for functional nanophotonics,” *Journal of Optics*, vol. 18, p. 103001, sep 2016.
- [56] T. Bucher, A. Vaskin, R. Mupparapu, F. J. F. Löchner, A. George, K. E. Chong, S. Fasold, C. Neumann, D.-Y. Choi, F. Eilenberger, F. Setzpfandt, Y. S. Kivshar, T. Pertsch, A. Turchanin, and I. Staude, “Tailoring photoluminescence from mos2 monolayers by mie-resonant metasurfaces,” *ACS Photonics*, vol. 6, no. 4, pp. 1002–1009, 2019.
- [57] Q. Le-Van, X. Le Roux, A. Aassime, and A. Degiron, “Electrically driven optical metamaterials,” *Nature Communications*, vol. 7, no. 12017, pp. 2041–1723, 2016.