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Efficient Bragg diffraction in thin semiconductor two-dimensional gratings

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Highly improved diffraction properties are demonstrated in a two-dimensional [2D] grating consisting of a transmission grating optically recorded in a semiconductor one-dimensional photonic crystal (1D-PC). Near unity internal diffraction efficiency, high wavelength selectivity, and Bragg diffraction regime operation are demonstrated when the read beam is set at Bragg incidence on the transmission grating while its wavelength corresponds to the band edge of the 3 μm thick 1D-PC. When the 2D grating is grown on a Bragg mirror, a single diffracted beam is obtained, which makes the device promising for optical signal processing.

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Optical signal processing greatly benefits from the use of transmission diffraction gratings [1]. For the purpose of compactness, thin nonlinear materials such as polymers, liquid crystals, or semiconductor nanostructures are used [2–5]. However, their small thicknesses imply a Raman–Nath diffraction regime with several diffraction orders, which limits the maximum power reached in the desired order. Such a drawback was circumvented by placing the holographic material in an asymmetric cavity [6–8], and the Bragg operation of the device was obtained, even with a thin recording material [9,10]. Recently, another solution using two-dimensional (2D) photonic crystals was proposed and demonstrated in photopolymers [11]. However, due to the low refractive index contrast ($\Delta n=0.045$) that could be reached in such materials, the device performances remained limited. In this Letter we show that the use of semiconductor 2D gratings exhibiting much larger refractive index contrasts in the mean direction of the read and diffracted beams enables Bragg internal diffraction efficiencies of unity together with high wavelength selectivity, even for micrometer thick devices.

As explained in detail in [11], the technique of band-edge-induced Bragg diffraction in 2D gratings is based on the enhancement of light intensity at the band edge wavelength of the one-dimensional (1D) refractive index modulation along the mean propagation direction \hat{Z} of the read and diffracted beams. Indeed, for a transmission grating with a refractive index modulation perpendicular to the \hat{Z} direction, very strong enhancement of the diffraction efficiency is predicted if the read beam propagation direction is set at Bragg incidence on this transmission grating while its wavelength is at the band edge resonance of the one-dimensional photonic crystal (1D-PC). In such a case, both read and -1 diffraction order intensities are enhanced through the local field amplification [12,13]. Higher diffraction orders are emitted at

different angles and only benefit from the read beam resonance. Their intensities can be negligible, especially in the case of low efficiency transmission gratings. Indeed, three advantages come from using such a 2D grating device: a huge enhancement of the diffraction efficiency (a factor of 10 to 1000), a largely improved wavelength selectivity (a factor of a few hundreds), and a Bragg diffraction regime with only two diffracted beams (one in transmission, the other one in reflexion). A single diffracted beam counter-propagating the read beam can be obtained if the 1D-PC is deposited on a Bragg mirror [14].

Such a device was realized and tested using the experimental setup shown in Fig. 1. The sample S with a total thickness of 2.44 μm consisted of a 1D-PC containing 30 pairs of CdMnTe–CdMgTe layers with respective refractive indices of 3.088 and 2.585 and a band edge wavelength around 694 nm, deposited on a Bragg mirror that was totally reflecting at this wavelength [15]. The refractive index transmission grating was recorded in the 1D-PC thanks to the generation of free carriers through two-photon absorption at a wavelength of 820 nm that is below the optical bandgap of the 1D-PC. The two write beams of equal intensities interfered on the sample after being sym-

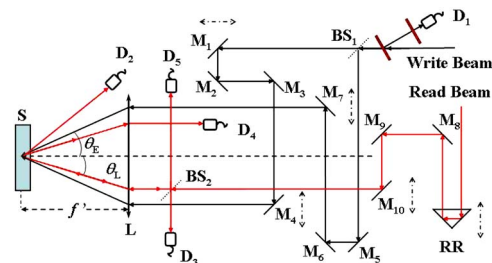


Fig. 1. (Color online) Experimental setup. Detector D_1 was used to monitor the pump energy. Detectors D_3 and D_4 , D_3 and D_5 , and D_3 and D_2 allowed the measurement of the signal reflectivity, -1 order diffraction efficiency, and $+1$ order diffraction efficiency, respectively.

metrically focused by a single lens. For a grating period of $4.46 \mu\text{m}$ the write incidence angle θ_W was adjusted to 5.3° thanks to the moving mirrors M_4 and M_7 , and the synchronization of the two write pulses was ensured by the moving mirror set M_1 – M_2 . The write beams came from one of the two outputs of a dual independently tunable-wavelength optical parametric generator (OPG) pumped by the same frequency-tripled mode locked Nd-YAG laser [16]. The second output of the OPG was tuned around 694 nm and used as the read beam. This read beam was focused on the sample at the Bragg incidence angle $\theta_R=4.5^\circ$ of the recorded transmission grating. The Bragg direction was adjusted thanks to the moving mirror M_{10} . The net diffraction efficiency (the measured intensity ratio of the reflected diffracted and incident read beams) was calibrated by replacing the sample by a gold mirror set orthogonally to the read beam.

A spectroscopic study was first undertaken for a write-read delay ($\Delta t \approx 20 \text{ ps}$) larger than the pulse duration to maximize the free carrier density generated by the interfering write beams. Figure 2 shows both the read beam reflectivity spectra recorded at a high write fluence (0.62 mJ/cm^2) and in the linear regime with no write beams [Fig. 2(a)] and the net diffraction efficiency spectra of this read beam for the -1 and $+1$ diffraction orders [Fig. 2(b)]. These spectra were recorded around the band edge wavelength and the hole at 694.1 nm in the linear reflectivity spectrum was due to diffusion losses at the interfaces of the CdMnTe–CdMgTe layers enhanced by the high local field at the band edge wavelength [13]. At a high write fluence the reflectivity spectrum is blueshifted [15] by 0.27 nm , which corresponds to a mean refractive index change $\Delta n = 1.2 \times 10^{-3}$ determined using the procedure described in [17]. For two write beams of equal intensities this gives rise to a refractive index modulation $n_1 = 3\Delta n_0/2 = 1.8 \times 10^{-3}$ (free carriers are generated by two-photon absorption), and the diffraction efficiency of the same transmission grating written in a homogeneous medium can be calculated from $\rho_{D \text{ hom}} = (\pi n_1 L/\lambda)^2 = 1.9 \times 10^{-4}$.

Figure 2(b) shows the huge enhancement of the Bragg (-1 order) diffraction at the band edge frequency. The diffraction efficiency increases from about 2.5×10^{-4} (very close to that estimated for the homogenous medium) far from the band edge resonance to around 1% at this band edge resonance (an increase by a factor of 40) despite losses introduced

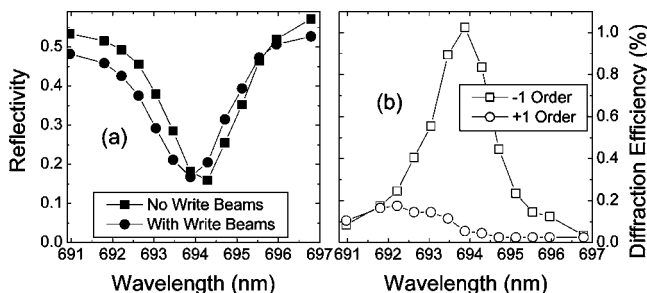


Fig. 2. (a) Reflectivity and (b) diffraction efficiency spectra.

by the nanostructuration of the holographic medium. Moreover the device is much more resonant ($\Delta\lambda_{CP} \approx 1.8 \text{ nm}$) than the same transmission grating recorded in a homogeneous medium ($\Delta\lambda_{CP} \approx 10,000 \text{ nm}$ calculated for this $1.9 \mu\text{m}$ thick grating). Finally, Fig. 2(b) also represents the net diffraction efficiency spectrum of the $+1$ diffraction order. It is much less intense than order -1 (a factor of 18 at the band edge) as its diffraction angle is 13.6° . Such behavior demonstrates that a Bragg diffraction regime can be obtained for a very thin 2D grating, which makes this device very attractive for the purpose of optical signal processing.

The maximum net diffraction efficiency obtained at the band edge of the 1D-PC was also measured as a function of the total pump fluence. Experimental results are shown in Fig. 3, together with a fit taking into account the two-photon absorption-induced free carrier origin of the refractive index modulation. The net diffraction efficiency ρ_{net} can then be phenomenologically described by

$$\rho_{\text{net}} = T_{\text{CP}}^2 \exp(-2\alpha F_p^2) \sin^2(\Delta\Phi F_p^2), \quad (1)$$

where $1 - T_{\text{CP}}$ and $1 - \exp(-\alpha F_p^2)$ represent optical losses for both read and diffracted beams, respectively, due to the linear diffusion at the CdMnTe–CdMgTe interfaces (the value $T_{\text{CP}} = 0.17$ was taken from Fig. 2) and the linear absorption due to the free carriers photogenerated through two-photon absorption of the write beams of total fluence F_p ($\alpha = 0.08 \text{ cm}^4 \text{ mJ}^{-2}$ was determined by the fit); $\rho_{\text{int}} = \sin^2(\Delta\Phi F_p^2)$ represents the internal diffraction efficiency of the device at the write fluence F_p ($\Delta\Phi = 0.27 \text{ cm}^4 \text{ mJ}^{-2}$ was determined by the fit). The excellent quality of the fit is actually a validation of the phenomenological Eq. (1). The fit also demonstrated that the internal conversion efficiency nearly reached 100% (93% at the maximum pump fluence used). The small value (1.3%) of the measured net diffraction efficiency was then only due to losses (a factor of about 35 for the diffusion losses and a factor of 2 for free carrier absorption losses). Considering the possibility of applying such a device to optical signal processing, these large losses indicate that the transmission grating had better be recorded either optically using electronic Kerr effect or, even better, electrically using Pockels effect with a comb electrode. More important, the linear losses should be almost suppressed by increasing the sample quality or (and) using thinner 1D-PCs of a larger refractive index contrast [17].

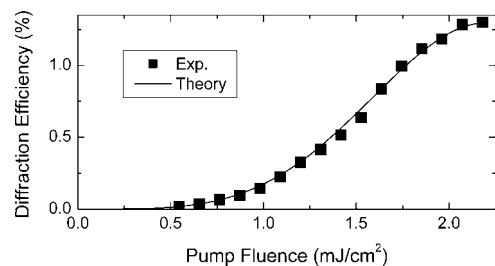


Fig. 3. Diffraction efficiency versus write fluence.

Time-resolved write-read experiments were also performed to verify the postulated origin of the refractive index modulation. Figure 4 shows the experimental net diffraction efficiency $\rho_{\text{net}}(\delta t)$ of the grating plotted as a function of the write-read delay δt for a pump fluence of 1.75 mJ cm^{-2} together with a fit performed using the expression [derived from Eq. (1)]

$$\rho_{\text{net}}(\delta t) = \frac{T_{\text{CP}}^2}{\sqrt{\pi\Delta t}} \int_{-\infty}^{+\infty} \sin^2[\Phi(t)] \exp[-2\alpha(t)] \times \exp\left[-\left(\frac{t-\delta t}{\Delta t}\right)^2\right] dt, \quad (2)$$

with $\Phi(t) = \sqrt{2/\pi}(\Phi_0/\Delta t) \int_{-\infty}^t \exp[-2\tau^2/\Delta t^2] d\tau$ and $\alpha(t) = \sqrt{2/\pi}(\alpha_0/\Delta t) \int_{-\infty}^t \exp[-2\tau^2/\Delta t^2] d\tau$, where $\Phi_0=0.82$ and $\alpha_0=0.23$ were given by the fit of the experimental results shown in Fig. 3. The pulse duration Δt was then the only free parameter in this fit. The very high quality of the fit shown in Fig. 4 for $\Delta t=10 \text{ ps}$ is further proof of the validity of Eq. (1). The long plateau observed at large delays shows that carrier diffusion is negligible in the results of Fig. 3, but it could also be detrimental to fast optical processing. From this point of view also, instantaneous electronic Kerr and Pockels effects would evidently be much more interesting for the recording of the transmission gratings.

The near unity internal conversion efficiency was obtained using a linearly polarized read beam. As in our device the diffracted beam is counterpropagating to the read beam; the best effective conversion efficiency (the intensity ratio of the diffracted beam coming out of the device to that of the read laser) is only 25% (when using a 50–50 beam splitter). As mentioned earlier, this drawback could in principle be circumvented by using a circularly polarized read beam, provided the response of the diffraction grating is isotropic [14]. This proposition was tested by replacing the beam splitter BS₂ by a polarizing beam splitter and a quarter-wave plate. The results were exactly the same, but the effective diffraction efficiency was measured rather than the net one. As expected all the diffracted light was collected at the output of the

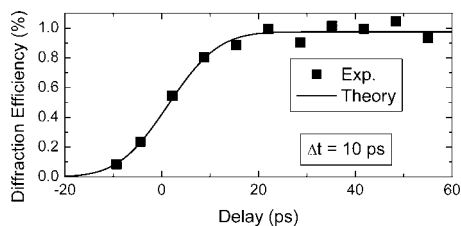


Fig. 4. Diffraction efficiency versus write-read delay.

polarizing beam splitter, which is very promising for application to dense optical signal processing.

In conclusion, diffraction of transmission gratings recorded orthogonally to a semiconductor 1D-PC was studied at the band edge wavelength of the 1D-PC. For this very thin 2D grating, an internal diffraction efficiency close to unity was measured for the transmission grating with an increase by a factor of 4,000 when compared to the same grating recorded in a homogeneous semiconductor. The device is also wavelength selective (2 nm) and operates in the Bragg diffraction regime with only one diffracted beam provided a suitable polarizing setup is used. This technique is very promising for dense optical signal processing provided optical losses of the 1D-PC are lowered, for instance, by using a very thin high-index contrast 1D-PC for the inscription of the transmission grating.

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